

Role of IR-Improvement in Precision LHC/FCC Physics and in Quantum Gravity

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IR-improvement based on amplitude-level resummation allows one to control unintegrable results in quantum field theory with arbitrary precision in principle. We illustrate such resummation in specific examples in precision LHC and FCC physics and in quantum gravity.

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1. Introduction

With the era of precision LHC physics squarely upon us, the relevance of resummation in quantum field theory is large. This relevance is amplified by the need to plan for future colliders in view of the current lack of experimental evidence for new degrees of freedom beyond those in the Standard Theory¹ in the LHC data at this writing. Indeed, several proposals for the next big colliding beam device involve a precision physics e^+e^- colliding beams phase in the regime of \sqrt{s} from 91 GeV to 390 GeV as a first step: FCC, ILC, CEPC, and CLIC. We mention that, in the problem of the union of quantum mechanics with gravity as given by the Einstein-Hilbert theory, precision theory methods seem also to be relevant. In what follows, we discuss the role of IR-improvement in the particular approach amplitude-based resummation in quantum field theory as it is applied to the aforementioned precision physics considerations.

In discussing the role of IR-improvement in quantum field theory in precision physics considerations, we are called to address the general idea of resummation in quantum field theory for IR, UV or CL (collinear) effects. The matter can be put into focus by comparing with the simple example of summation for a geometric series, such as

$$\frac{1}{1-x} = \sum_{n=0}^{\infty} x^n. \quad (1.1)$$

Here the result of the summation is an analytic result that is well defined except for a simple pole at $x = 1$. While the mathematical tests for convergence on the simple series would only guarantee convergence for $|x| < 1$, the result of the summation yields a result that is well-behaved in the entire complex plane except for a simple pole at $x = 1$. The message is that infinite order summation can yield results that are much better behaved than what one can deduce order-by-order in the respective series. One is accordingly invited, as a next step, to resum series that are already being summed to try to improve ones knowledge of the represented function.

Continuing in this way, we write for the Feynman series for a process the following:

$$\sum_{n=0}^{\infty} C_n \alpha_s^n \left\{ \begin{array}{l} = F_{\text{RES}}(\alpha_s) \sum_{n=0}^{\infty} B_n \alpha_s^n, \text{ EXACT} \\ \cong G_{\text{RES}}(\alpha_s) \sum_{n=0}^{\infty} \tilde{B}_n \alpha_s^n, \text{ APPROX.} \end{array} \right. \quad (1.2)$$

Two versions of resumming the original Feynman series for the quantity on the LHS are given on the RHS. One is an exact re-arrangement of the original series. It is labeled ‘exact’. The other only agrees with the LHS to some fixed order N in the expansion parameter α_s . It is labeled ‘approx’. Which version of resummation is to be preferred has been discussed for some time [2]. At this writing, the discussion continues to new paradigms.

In QCD \otimes QED the discussion continues with added dimensions. We have the ongoing comparison, by two of us (BFLW and SAY), of the hard IR cut-offs in MC’s such as Herwig [3], Herwig++ [4], Pythia [5] and Sherpa [6] versus the resummed integrability of the IR in Herwig1.031 [7]. In addition, we have the approaches to ISR radiation from quarks in hard hadron-hadron scattering processes, wherein we have the approach with QED PDF’s [8] versus exact Feynman diagrams with short-distance quark masses as pursued by four of us (SJ, BFLW, ZAW, and

¹We follow Prof. D.J. Gross [1] and henceforth refer to the Standard Model as the Standard Theory of elementary particles.

SAY) in Ref. [9]. In the FCC and other future collider frameworks, the attendant precision physics programs necessitate IR-improved quantum field theory results to exploit properly the respective new physics windows. In quantum gravity, it has been found that resummed IR leads to UV finiteness [10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25]. In what follows, we expand on these developments in IR-improvement of quantum field theory from the standpoint of applications to precision LHC/FCC physics and applications of one of us (BFLW) of quantum gravity.

We close this introductory section by calling attention to the 50-year celebrations of the Standard Model of Elementary Particles [26, 27, 28, 29, 30, 31, 32, 33]. At the same time, plans are underway for the future explorations in our field as we have noted. We call attention here to the FCC Study [34] which we show in Fig. 1. The plans call for a precision studies e^+e^- phase featur-

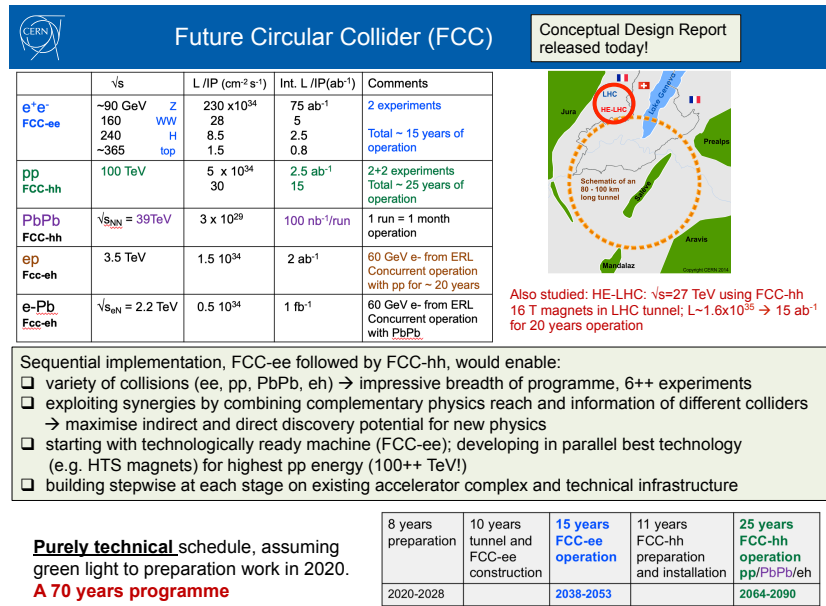


Figure 1: Tentative schedule and physical layout in FCC Study as portrayed in Ref. [34].

ing more than 5 Tera Z's, for example, followed by a 100 TeV hadron collider, reminiscent of the canceled SSC in Waxahachie, TX, USA. Only now do we really see how big of a mistake the latter cancellation was. It is important to keep this historical perspective in mind.

The discussion proceeds as follows. In the next section, we review briefly exact amplitude-based resummation theory. This is followed by a discussion of applications in precision LHC physics in Section 3. Applications in Precision FCC physics are discussed in Section 4. Section 5 present applications in quantum gravity. Section 6 contains our summary remarks.

2. Review of Exact Amplitude-Based Resummation Theory

This approach to higher order corrections is founded on the master formula

$$d\bar{\sigma}_{\text{res}} = e^{\text{SUM}_{\text{IR}}(\text{QCED})} \sum_{n,m=0}^{\infty} \frac{1}{n!m!} \int \prod_{j_1=1}^n \frac{d^3 k_{j_1}}{k_{j_1}}$$

$$\prod_{j_2=1}^m \frac{d^3 k'_{j_2}}{k'_{j_2}} \int \frac{d^4 y}{(2\pi)^4} e^{iy \cdot (p_1 + q_1 - p_2 - q_2 - \sum k_{j_1} - \sum k'_{j_2}) + D_{\text{QCD}}} \tilde{\beta}_{n,m}(k_1, \dots, k_n; k'_1, \dots, k'_m) \frac{d^3 p_2}{p_2^0} \frac{d^3 q_2}{q_2^0}, \quad (2.1)$$

where *new* (YFS-style) *non-Abelian* residuals $\tilde{\beta}_{n,m}(k_1, \dots, k_n; k'_1, \dots, k'_m)$ have n hard gluons and m hard photons. The infrared functions $\text{SUM}_{\text{IR}}(\text{QCD})$ and D_{QCD} and the residuals are defined in Ref. [35, 36]. In the implementation of parton shower/ME matching, we have the replacements $\tilde{\beta}_{n,m} \rightarrow \hat{\beta}_{n,m}$. These replacements allow us, via the basic formula

$$d\sigma = \sum_{i,j} \int dx_1 dx_2 F_i(x_1) F_j(x_2) d\hat{\sigma}_{\text{res}}(x_1 x_2 s), \quad (2.2)$$

to connect with MC@NLO [37, 38] as explained in Ref. [35, 36].

We have made various applications of Eq.(2.1) to precision LHC and FCC physics and we have made applications of its extension to general relativity as an approach to quantum gravity. We discuss such applications in precision LHC physics in the next Section.

3. Applications in Precision LHC Physics

In Refs. [35, 36, 39], two of us (BFLW and SAY) have interfaced the realization of IR-improved DGLAP-CS [40, 41, 42, 43, 44, 45, 46, 47, 48] theory as realized in Herwig1.031 [7] with MC@NLO and with MG5_aMC@NLO [49] for precision studies of Z and W+jets production at the LHC. In Refs. [9, 50, 51] four of us (SJ, BFLW, ZAW, and SAY) have realized exact $O(\alpha^2 L)$ CEEX EW corrections for hadron-hadron scattering processes in the MC $\mathcal{K}\mathcal{K}\text{MC-hh}$ which features built-in parton showers from Herwig6.5 and Herwig1.031 as well as an LHE [52] format to facilitate the interface with other parton shower MC's. In Z and W+ jets productions, IR-improvement gives a comparable or better data fit without ad hoc parameters for the intrinsic transverse momentum spectrum of partons inside the proton. In $\mathcal{K}\mathcal{K}\text{MC-hh}$, IR-improvement allows to quantify role of ISR in precision predictions for Z production observables, as we now illustrate.

Specifically, our implementation for DIZET [53, 54] in $\mathcal{K}\mathcal{K}\text{MC-hh}$ uses a modified G_μ scheme with Standard Theory² input parameters as given in Ref. [55] with the exception that we use the light quark masses $m_u = 6.0$ MeV, $m_d = 10.0$ MeV ($m_u = 2.2$ MeV, $m_d = 4.7$ MeV) where the latter set is used to estimate the uncertainty of our results due to the uncertainty in the precise values of the respective short-distance quark masses. For a systematic discussion of the effects of ISR and IFI as predicted by $\mathcal{K}\mathcal{K}\text{MC-hh}$ on the angular observables in Z production and decay to lepton pairs at the LHC see Ref. [56]. Here, we focus on the use of the Z production observables to quantify the systematics in the world-leading ATLAS measurement [57] of M_W :

$$M_W = 80370 \pm 7(\text{stat.}) \pm 11(\text{exp. syst.}) \pm 14(\text{mod. syst.}) \text{ MeV} = 80370 \pm 19 \text{ MeV},$$

where the first uncertainty is statistical, the second is the experimental systematic uncertainty, and the third is the physics-modeling systematic uncertainty. Z/γ^* production with decay to lepton pairs are used in determining the latter uncertainty. Specifically, the distributions for the lepton p_T spectra in single Z/γ^* in Ref. [57], as reproduced here in Fig. 2, are used to analyze expectations for the corresponding lepton spectra in the single W production. In both spectra featured in Fig. 2, we see

²We follow D.J. Gross [1] and henceforth refer to the Standard Model [26, 27, 28, 29, 30, 31] as the Standard Theory.

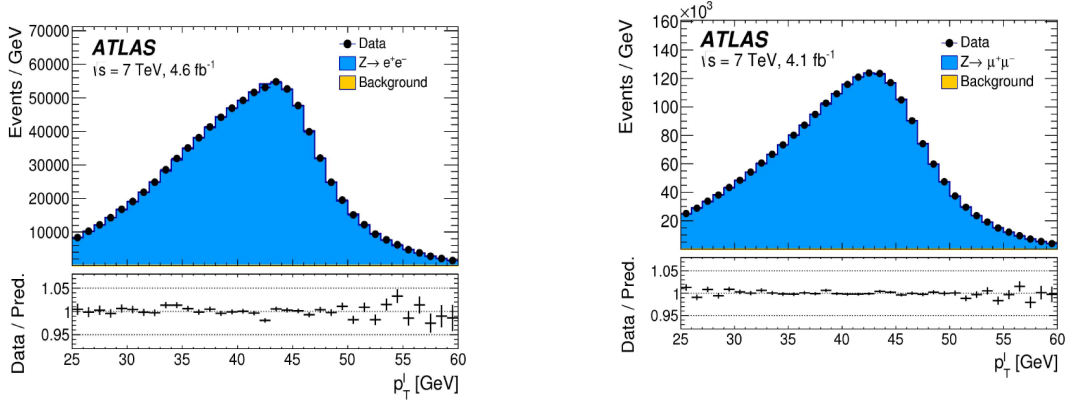


Figure 2: ATLAS electron and muon p_T spectra from Fig. 15 in Ref. [57].

that, at the lower end, the theory used by ATLAS is about 2% below the data and, at the high end, it is a similar amount above the data. We address these spectra in Ref. [50] as shown here in Fig. 3. In Fig. 3, the best prediction of the exact $O(\alpha^2 L)$ CEEX calculation (labeled CEEX2) is compared

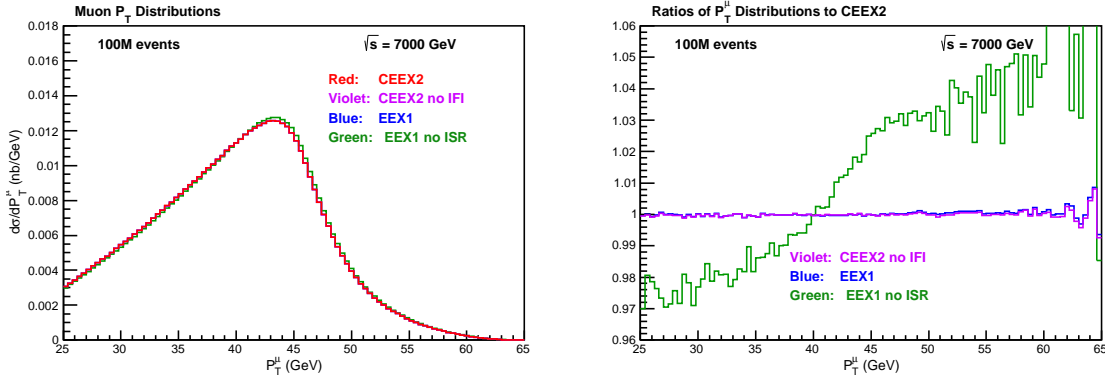


Figure 3: Muon transverse momentum distributions and their ratios for $\mathcal{K}\mathcal{K}MC$ -hh with the cuts specified in the text for the EW-CORR (electroweak-correction) labels “CEEX2” (red – medium dark shade), “CEEX2 (no IFI)” (violet – light dark shade), “EEX1” (blue – dark shade), and “EEX1 (no IFI)” (green – light shade), showered by HERWIG 6.5. The labels are explained in the text. The ratio plot features “CEEX2” as the reference distribution as noted in the respective title.

with less precise predictions denoted as follows: $O(\alpha^2 L)$ CEEX with ISR+FSR+IFI – labeled as “CEEX2”; $O(\alpha^2 L)$ CEEX without IFI (initial state final state interference) – labeled “CEEX2 (no IFI)”; $O(\alpha)$ EEX – labeled “EEX1”; $O(\alpha)$ EEX without ISR (initial state radiation) – labeled “EEX1 (no ISR)”. Here, EEX denotes YFS exponentiation at the level of squared amplitudes as originally formulated in Ref. [58] in contrast to CEEX, which represents the coherent realization of YFS exponentiation at the level of the amplitudes, as explained in Ref. [59]. We see that the EEX1 result with no ISR is about 2% below the best calculation on the low end of the spectrum and is a similar amount above the best calculation on the high end, just as we see in the ATLAS data, where the theory prediction also has no QED ISR other than that in its Pythia [60] shower, which is essentially collinear by construction. We conclude that the corrections in $\mathcal{K}\mathcal{K}MC$ -hh should be

implemented into the ATLAS analysis and other similar analyses.

4. Applications in Precision FCC Physics

The FCC project involves a 100 km tunnel in which one first installs an e^+e^- collider that operates for precision Z, WW, Higgs and top physics, to be followed by the installation of a 100 TeV proton-proton collider. The first phase is FCC-ee, the second phase is FCC-hh. As three of us (BFLW, ZAW and SAY) have argued in Ref. [51], IR-improvement of even the FCC-hh discovery spectra is needed. In the FCC-ee, more than 5 Tera Z's will be observed³, so that a key issue is the theoretical precision of the luminosity as discussed by five of us (SJ, WP, MS, BFLW, SAY) in Ref. [66], a discussion to which we now turn.

To set the perspective, we show in Tab. 1 the state of the art at the time of LEP, where the precision tag on the theoretical prediction of the luminosity was 0.061%(0.054%), according to

Type of correction/error	LEP1		LEP2	
	1996	1999	1996	1999
(a) Missing photonic $O(\alpha^2)$ [67, 68]	0.10%	0.027%	0.20%	0.04%
(b) Missing photonic $O(\alpha^3 L_e^3)$ [69]	0.015%	0.015%	0.03%	0.03%
(c) Vacuum polarization [70, 71]	0.04%	0.04%	0.10%	0.10%
(d) Light pairs [72, 73]	0.03%	0.03%	0.05%	0.05%
(e) Z and s-channel γ [74, 75]	0.015%	0.015%	0.0%	0.0%
Total	0.11% [75]	0.061% [76]	0.25% [75]	0.12% [76]

Table 1: Summary of the total (physical+technical) theoretical uncertainty for a typical calorimetric detector. For LEP1, the above estimate is valid for a generic angular range within 1° – 3° (18–52 mrad), and for LEP2 energies up to 176 GeV and an angular range within 3° – 6° . Total uncertainty is taken in quadrature. Technical precision included in (a).

whether one does not(does) treat the effects of soft pairs [77, 78]. See Ref. [76] for a discussion of the various error estimates as they relate to the cited published works in the table. The current state of the art, 0.038%, is shown in Table 2, where we note that recently this precision tag has been effected in Ref. [83] in reducing the error on the number, N_ν , of light neutrinos so that it is now $N_\nu = 2.9975 \pm 0.00742$. The FCC-ee requirement [84] on the theoretical precision of the luminosity is 0.01%. The steps involved in realizing this requirement are discussed in Ref. [66]. The conclusion of Ref. [66] is that, with sufficient resources to make the required improvements to the MC `BHLUMI` [85], the 0.01% precision tag can be realized, as we show in Table 3. Again, we cite the published works associated to the error estimates in the table as discussed in Ref. [66]. The current planning for the FCC [84, 86] includes a provision for such resources as those called for in Ref. [66].

³For this scenario, recently, thanks to the modular program design improvements for the electroweak sector calculations became available in `KKMC` [59]. The comparison of numerical results is assured. This is of importance for future long-term projects as discussed in [61]. In particular the up-to-date parametrization of $\alpha_{QED,eff}$ [62] (section B.1) is available as well as the electroweak library `DIZET` version 6.42 [63] and the recent version 6.45. The same is true and is of importance for precision measurements of electroweak effects with τ lepton polarization [64] where benefits from future Belle experimental results on τ lepton decays are of a sizable potential. In this context technical tests embedded in the older version of `TAUOLA` for the τ lepton decay [65] may need to be revisited to assure technical correctness at better than the precision standard of today.

Type of correction / Error	1999	Update 2018
(a) Photonic $O(L_e\alpha^2)$	0.027% [68]	0.027%
(b) Photonic $O(L_e^3\alpha^3)$	0.015% [69]	0.015%
(c) Vacuum polariz.	0.040% [70, 71]	0.013% [79]
(d) Light pairs	0.030% [73]	0.010% [77, 78]
(e) Z and s-channel γ exchange	0.015% [80, 81]	0.015%
(f) Up-down interference	0.0014% [82]	0.0014%
(f) Technical Precision	–	(0.027)%
Total	0.061% [76]	0.038%

Table 2: Summary of the total (physical+technical) theoretical uncertainty for a typical calorimetric LEP luminosity detector within the generic angular range of 18–52 mrad. Total error is summed in quadrature.

Type of correction / Error	Update 2018	FCC-ee forecast
(a) Photonic [$O(L_e\alpha^2)$] $O(L_e^2\alpha^3)$	0.027%	0.1×10^{-4}
(b) Photonic [$O(L_e^3\alpha^3)$] $O(L_e^4\alpha^4)$	0.015%	0.6×10^{-5}
(c) Vacuum polariz.	0.014% [79]	0.6×10^{-4}
(d) Light pairs	0.010% [77, 78]	0.5×10^{-4}
(e) Z and s-channel γ exchange	0.090% [74]	0.1×10^{-4}
(f) Up-down interference	0.009% [82]	0.1×10^{-4}
(f) Technical Precision	(0.027)%	0.1×10^{-4}
Total	0.097%	1.0×10^{-4}

Table 3: Anticipated total (physical+technical) theoretical uncertainty for a FCC-ee luminosity calorimetric detector with the angular range being 64–86 mrad (narrow), near the Z peak. Description of photonic corrections in square brackets is related to the 2nd column. The total error is summed in quadrature.

5. Applications in Quantum Gravity

We turn now to the role of IR-improvement in quantum gravity. The question of whether or not quantum gravity is calculable in relativistic quantum field theory is still open in the literature, as one of us (BFLW) has discussed in Refs. [87, 88]. Here, one of us argues that it is if he extends the YFS [58, 89, 59] version⁴ of the exact resummation example to resum the Feynman series for the Einstein-Hilbert Lagrangian for quantum gravity. The resultant resummed theory, resummed quantum gravity (RQG), in analogy with what we see in the example in Eq.(1.1), is very much better behaved in the UV compared to what one would estimate from that Feynman series.

One of us (BFLW) has shown [22, 23, 24, 25] that the RQG realization of quantum gravity leads to Weinberg’s [91] UV-fixed-point behavior for the dimensionless gravitational and cosmological constants and that the resummed theory is actually UV finite. He has shown further that the RQG theory, taken together with the Planck scale inflationary [92, 93] cosmology formulation in Refs. [94, 95]⁵ from the asymptotic safety approach to quantum gravity in Refs. [10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21], allows us to predict the cosmological constant Λ . Specifically, he obtains, employing the arguments in Refs. [97] (t_{eq} is the time of matter-radiation equality), the

⁴YFS-type soft resummation and its extension to quantum gravity was also worked-out by Weinberg in Ref. [90].

⁵The authors in Ref. [96] also proposed the attendant choice of the scale $k \sim 1/t$ used in Refs. [94, 95].

result [98]

$$\begin{aligned}\rho_\Lambda(t_0) &\cong \frac{-M_{Pl}^4(1 + c_{2,eff}k_{tr}^2/(360\pi M_{Pl}^2))^2}{64} \sum_j \frac{(-1)^F n_j}{\rho_j^2} \times \frac{t_{tr}^2}{t_{eq}^2} \times \left(\frac{t_{eq}^{2/3}}{t_0^{2/3}}\right)^3 \\ &\cong \frac{-M_{Pl}^2(1.0362)^2(-9.194 \times 10^{-3})(25)^2}{64} \frac{(25)^2}{t_0^2} \cong (2.4 \times 10^{-3} eV)^4.\end{aligned}\quad (5.1)$$

$t_0 \cong 13.7 \times 10^9$ yrs is the age of the universe, $t_{tr} \sim 25t_{Pl}$ is the transition time between the Planck regime and the classical Friedmann-Robertson-Walker(FRW) regime in the Planck scale cosmology description of inflation in Ref. [95], $c_{2,eff} \cong 2.56 \times 10^4$ is defined in Refs. [22, 23, 24], t_{eq} is the time of matter-radiation equality, and M_{Pl} is the Planck mass.

Due to the prediction's closeness to the observed value [99, 100, 101], $\rho_\Lambda(t_0)|_{\text{expt}} \cong ((2.37 \pm 0.05) \times 10^{-3} eV)^4$, one of us (BFLW) discusses in Ref. [102] its reliability and argues that its uncertainty is at the level of a factor of $O(10)$. There follow constraints on susy GUT's as well as discussed in Refs. [98]. Various consistency checks on our result for $\rho_\Lambda(t_0)$ are also considered in Ref. [98].

6. Summary

What we can see is that the resummation of the IR regime of quantum field theory, coupled with exact results to a given order, when done at the level of the amplitude in an exact rearrangement of the original Feynman series, has wide applicability for precision phenomenology. We have illustrated the range of this applicability from the current precision LHC physics to the futuristic FCC precision physics program and reliable estimates for quantum gravity by one of us (BFLW). The future of particle physics is intimately interwoven with continued progress on IR-improved quantum field theoretic predictions, when taken together with control of the attendant collinear and UV limits.

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